

# Calculation of Energy Levels and B (E2) for $^{46-45}\text{Ca}$ , $^{46-45}\text{Sc}$ and $^{46-45}\text{Ti}$ by Using Nuclear Shell Model Code OXBASH

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**Abstract:** In this paper, the energy levels for  $^{46-45}\text{Ca}$ ,  $^{46-45}\text{Ti}$  and  $^{46-45}\text{Sc}$  isotopes have been calculated. Calculations were carried out in the f7-shell region by employing the effective interactions, f748pn and using the shell model code OXBASH for windows by applying spin-parity of valance nucleons. It is found that there is good convergence of energy levels values with the standard practical value.

**Key words:** Energy levels, OXBASH Code and gamma transitions.

## 1. Introduction

The aim of this paper is study of the energy levels of  $^{46-45}\text{Ca}$ ,  $^{46-45}\text{Ti}$  and  $^{46-45}\text{Sc}$  isotopes by using OXBASH code for windows. This program is a set of codes carrying out shell-model calculations with dimensions up to about 50,000 in the J-T scheme and about 2,000,000 in the M-scheme. Oxbash comes with a library of model spaces and interactions [1, 2].

Applied the shell model and use a Modified Brown and Sherr (f748pn) interaction for neutron and proton orbits in  $^{46-45}\text{Ca}$ ,  $^{46-45}\text{Ti}$  and  $^{46-45}\text{Sc}$  to calculate the energy levels values. Various observables can be predicted accurately and systematically in terms of the nuclear shell model. For light nuclei, there are several "standard" effective interactions such as the Cohen-Kurath [3] and the USD [4] interactions for the *p* and *sd* shells, respectively. On the other hand, in the next major shell, *i.e.*, in the f7-shell, there are also "standard" interaction such as f748pn [5]. The spectroscopy of nuclei, in the f7-shell region, has been well described within the shell model framework. The best example for using several model spaces and two-body interactions is that of Brown et al, which is the most remarkable work in this field [6,7]. The starting point in all such shell-model calculations is the derivation of an effective interaction owing to the fact that the f7-shell is the most important for a variety of problems in nuclear structure such as electron capture in supernova explosions. In this work, the shell model calculations are carried out in the f7-shell region for the isotopes  $^{46-45}\text{Ca}$ ,  $^{46-45}\text{Ti}$  and  $^{46-45}\text{Sc}$ , to test the ability of the present effective interactions in reproducing the experiment in this mass region.

## 2. Shell Model Calculations

The calculations have been carried out in the nuclear shell model f7 using the code OXBASH for windows [6]. The code uses an m-scheme Slater determinant basis. With a projection technique, there are been constructed wave functions with good angular momentum J and isospin T. The f7pn model space is comprised of (1f7/2) below the closed N = Z=20 shell [8]. One can find the calculated results of states of the odd A and even A nuclei, number of protons, *i.e.*, 20 to the  $^{46-45}\text{Ca}$ , with neutron numbers (25, 26), number of protons, 21 to the  $^{46-45}\text{Sc}$ , with neutron numbers (24, 25) and number of protons, 22 to the  $^{46-45}\text{Ti}$ , with neutron numbers (23, 24) energy levels value.

### 2.1. Energy Levels Calculations

The calculations have been carried out using the code OXBASH for windows [9]. In the f7 model space comprised of the 0f7/2 valence orbits outside the <sup>40</sup>Ca. Two effective interactions have been employed with f7 model space for the calculations of level spectra, these effective interaction is f748pn [5]. It is worth mentioning that <sup>45</sup>Ca and <sup>46</sup>Ca have Isospin part (*T* = 2.5 and 3) respectively, <sup>45</sup>Sc and <sup>46</sup>Sc have Isospin part (*T* = 1.5 and 2) respectively, while <sup>45</sup>Ti and <sup>46</sup>Ti have Isospin part (*T* = 0.5 and 1) respectively.

The energy levels values for <sup>45</sup>Ca nucleus, the energy levels values are shown in table 1. these values are agreement with the experimental values, and new energy levels have been reached.

**Table 1** shows a comparison of the energy levels values with respect to the ground state were calculated from f748pn effective interactions with experimental excitation energies of <sup>45</sup>Ca

Exp .Res[10]		F748pn	<i>J<sup>π</sup></i>
<i>J<sup>π</sup></i>	Energy elevels		
7/2 <sup>-</sup>	0.000	0.000	7/2 <sub>1</sub> <sup>-</sup>
5/2 <sup>-</sup>	0.174	0.151	5/2 <sub>1</sub> <sup>-</sup>
3/2 <sup>-</sup>	1.434	1.221	3/2 <sub>1</sub> <sup>-</sup>
-----	1.584	1.616	11/2 <sub>1</sub> <sup>-</sup>
-----	1.94	1.927	9/2 <sub>1</sub> <sup>-</sup>
-----	2.953	2.936	15/2 <sub>1</sub> <sup>-</sup>

The energy levels values for <sup>46</sup>Ca nucleus from f748pn effective interaction is shown in table 2 and these effective interactions give good results in comparison with the experimental values.

**Table 2** shows a comparison of the energy levels values with respect to the ground state were calculated from f748pn effective interactions with experimental excitation energies of <sup>46</sup>Ca

Exp .Res[11]		F748pn	<i>J<sup>π</sup></i>
<i>J<sup>π</sup></i>	Energy elevels		
0 <sup>+</sup>	0.000	0.000	0 <sub>1</sub> <sup>+</sup>
2 <sup>+</sup>	1.346	1.346	2 <sub>1</sub> <sup>+</sup>
4 <sup>+</sup>	2.574	2.575	4 <sub>1</sub> <sup>+</sup>
6 <sup>+</sup>	2.973	2.974	6 <sub>1</sub> <sup>+</sup>

The energy levels values for <sup>45</sup>Sc nucleus from f748pn effective interaction is shown in table 3 and these effective interaction results reasonably consistent with experimental data. The total angular momentum and parity are ( 7/2<sub>1</sub><sup>-</sup>, 11/2<sub>1</sub><sup>-</sup>, 9/2<sub>1</sub><sup>-</sup>, 15/2<sub>1</sub><sup>-</sup>, 7/2<sub>3</sub><sup>+</sup>, 5/2<sub>2</sub><sup>-</sup>, 9/2<sub>3</sub><sup>-</sup>, 19/2<sub>1</sub><sup>-</sup>, 23/2<sub>1</sub><sup>-</sup>, 7/2<sub>7</sub><sup>-</sup>) respectively, confirmation of which is ( 7/2<sup>-</sup>, 15/2<sup>-</sup>, 21/2<sup>-</sup>, 5/2<sup>-</sup>) as well as confirmation of momentum only, which is (1/2, 3/2, 9/2, 5/2, 3/2) respectively, and new energy levels have been reached.

**Table 3** shows a comparison of the energy levels values with respect to the ground state were calculated from f748pn effective interactions with experimental excitation energies of <sup>45</sup>Sc

Exp .Res[10]		F748pn	<i>J<sup>π</sup></i>
<i>J<sup>π</sup></i>	Energy elevels		
7/2 <sup>-</sup>	0.000	0.000	7/2 <sub>1</sub> <sup>-</sup>
11/2 <sup>-</sup>	1.236	1.268	11/2 <sub>1</sub> <sup>-</sup>

$9/2^-$	1.662	1.436	$9/2_1^-$
-----	1.716	1.707	$3/2_1^-$
-----	1.900	1.883	$5/2_1^-$
$15/2^-$	2.090	1.969	$15/2_1^-$
$(1/2, 3/2, 5/2)$	2.151	2.298	$1/2_1^-$
-----	2.152	2.416	$13/2_1^-$
$(7/2^-)$	2.341	2.461	$7/2_2^-$
$(1/2^+, 3/2, 5/2)$	2.531	2.490	$3/2_2^-$
-----	2.700	2.730	$11/2_2^-$
$(9/2^+, 11/2^-)$	2.960	2.838	$9/2_2^-$
$5/2^-, 7/2^-$	2.747	2.920	$7/2_3^-$
$1/2^+, 3/2, 5/2$	3.092	2.958	$5/2_2^-$
$5/2^-, 7/2^-, 9/2^-$	3.136	3.171	$9/2_3^-$
-----	3.198	3.181	$17/2_1^-$
-----	3.224	3.280	$13/2_2^-$
$(15/2^-)$	3.363	3.307	$15/2_2^-$
+	3.443	3.424	$11/2_3^-$
$19/2^-$	3.692	3.556	$19/2_1^-$
-----	3.722	3.691	$15/2_3^-$
-----	4.129	4.131	$9/2_4^-$
-----	4.178	4.190	$7/2_4^-$
-----	4.244	4.220	$5/2_3^-$
-----	4.307	4.260	$11/2_4^-$
-----	4.424	4.343	$13/2_3^-$
-----	4.464	4.448	$5/2_4^-$
-----	4.546	4.526	$17/2_2^-$
$(1/2, 3/2)$	4.662	4.565	$3/2_3^-$
-----	4.713	4.631	$11/2_5^-$
-----	4.713	4.667	$7/2_5^-$
-----	4.774	4.759	$13/2_4^-$
-----	4.801	4.792	$1/2_2^-$
-----	5.219	5.206	$15/2_4^-$
-----	5.261	5.272	$3/2_4^-$
-----	5.374	5.313	$9/2_5^-$
-----	5.374	5.402	$5/2_5^-$
$23/2^-$	5.418	5.563	$23/2_1^-$
$(21/2^-)$	5.710	5.672	$21/2_1^-$
-----	5.774	5.674	$7/2_6^-$
-----	5.834	5.843	$19/2_2^-$
-----	6.244	6.239	$11/2_6^-$
-----	6.332	6.347	$9/2_6^-$
$5/2^-, 7/2^-, 9/2^-$	6.551	6.419	$7/2_7^-$
-----	6.609	6.510	$17/2_3^-$
$(5/2^-, 7/2^-, 9/2^-)$	6.667	6.653	$5/2_6^-$
-----	-----	6.757	$13/2_5^-$
-----	7.696	7.618	$3/2_5^-$

-----	8.118	7.973	11/2 <sub>7</sub> <sup>-</sup>
-----	8.305	8.250	9/2 <sub>7</sub> <sup>-</sup>
-----	9.164	9.085	15/2 <sub>5</sub> <sup>-</sup>

For <sup>46</sup>Sc nucleus the energy levels calculations from f748pn effective interaction is shown in table 4. The effective interactions give results reasonably consistent with experimental data. The total angular momentum and parity are (3<sub>1</sub><sup>+</sup>, 5<sub>1</sub><sup>+</sup>, 2<sub>1</sub><sup>+</sup>, 5<sub>2</sub><sup>+</sup>, 4<sub>1</sub><sup>+</sup>, 7<sub>1</sub><sup>+</sup>, 3<sub>4</sub><sup>+</sup>, 1<sub>2</sub><sup>+</sup>, 4<sub>4</sub><sup>+</sup>, 0<sub>1</sub><sup>+</sup>) respectively, confirmation of which is (2<sup>+</sup>, 4<sup>+</sup>) respectively, as well as confirmation of momentum only is (3, 3) respectively, and new energy levels have been reached.

**Table 4** shows a comparison of the energy levels values with respect to the ground state were calculated from f748pn effective interactions with experimental excitation energies of <sup>46</sup>Sc

Exp .Res[11]		F748pn	J <sup>π</sup>
J <sup>π</sup>	Energy elevels		
6+	0.520	0.000	6 <sub>1</sub> <sup>+</sup>
4+	0.000	0.122	4 <sub>1</sub> <sup>+</sup>
3+	0.227	0.214	3 <sub>1</sub> <sup>+</sup>
5+	0.280	0.238	5 <sub>1</sub> <sup>+</sup>
2+	0.444	0.497	2 <sub>1</sub> <sup>+</sup>
5+	0.774	0.814	5 <sub>2</sub> <sup>+</sup>
4+	0.835	0.909	4 <sub>2</sub> <sup>+</sup>
7+	0.978	0.979	7 <sub>1</sub> <sup>+</sup>
-----	1.298	1.473	1 <sub>1</sub> <sup>+</sup>
3+,4+,(2+)	1.427	1.556	2 <sub>2</sub> <sup>+</sup>
-----	1.676	1.671	8 <sub>1</sub> <sup>+</sup>
-----	1.825	1.894	6 <sub>2</sub> <sup>+</sup>
(3)+	2.070	2.045	3 <sub>2</sub> <sup>+</sup>
-----	2.084	2.109	5 <sub>3</sub> <sup>+</sup>
-----	2.395	2.406	7 <sub>2</sub> <sup>+</sup>
(4+,5+)	2.431	2.409	4 <sub>3</sub> <sup>+</sup>
-----	2.486	2.472	9 <sub>1</sub> <sup>+</sup>
-----	2.534	2.583	2 <sub>3</sub> <sup>+</sup>
3+,4+	2.568	2.596	3 <sub>3</sub> <sup>+</sup>
3+	2.705	2.748	3 <sub>4</sub> <sup>+</sup>
1+	2.815	2.770	1 <sub>2</sub> <sup>+</sup>
-----	2.890	2.888	5 <sub>4</sub> <sup>+</sup>
-----	3.081	3.082	8 <sub>2</sub> <sup>+</sup>
-----	3.081	3.102	6 <sub>3</sub> <sup>+</sup>
4+	3.176	3.169	4 <sub>4</sub> <sup>+</sup>
-----	3.338	3.351	7 <sub>3</sub> <sup>+</sup>
-----	3.381	3.355	1 <sub>3</sub> <sup>+</sup>
-----	3.813	3.813	9 <sub>2</sub> <sup>+</sup>
-----	3.868	3.823	6 <sub>4</sub> <sup>+</sup>
-----	3.937	3.915	10 <sub>1</sub> <sup>+</sup>
-----	3.945	3.934	2 <sub>4</sub> <sup>+</sup>
-----	3.961	3.967	8 <sub>3</sub> <sup>+</sup>
-----	4.039	4.050	5 <sub>5</sub> <sup>+</sup>
-----	4.200	4.197	11 <sub>1</sub> <sup>+</sup>
-----	4.447	4.459	3 <sub>5</sub> <sup>+</sup>

-----	4.846	4.840	7 <sub>4</sub> <sup>+</sup>
-----	4.972	4.995	4 <sub>5</sub> <sup>+</sup>
0+	5.022	5.004	0 <sub>1</sub> <sup>+</sup>
-----	5.327	5.331	5 <sub>6</sub> <sup>+</sup>
-----	6.362	6.354	2 <sub>5</sub> <sup>+</sup>
-----	-----	7.470	4 <sub>6</sub> <sup>+</sup>
-----	-----	7.808	6 <sub>5</sub> <sup>+</sup>

The energy levels values for <sup>45</sup>Ti nucleus from f748pn effective interaction is shown in table 5 and these effective interaction results reasonably consistent with experimental data. The total angular momentum and parity are ( 9/2<sub>1</sub><sup>-</sup>, 11/2<sub>1</sub><sup>-</sup>, 7/2<sub>2</sub><sup>-</sup>, 15/2<sub>1</sub><sup>-</sup>, 17/2<sub>1</sub><sup>+</sup>, 5/2<sub>2</sub><sup>-</sup>, 23/2<sub>1</sub><sup>-</sup>, 27/2<sub>1</sub><sup>-</sup>) respectively, confirmation of which is ( 3/2<sup>-</sup>) as well as confirmation of momentum only, which is (3/2, 11/2, 13/2, 7/2, 17/2) respectively, and new energy levels have been reached.

**Table 5** shows a comparison of the energy levels values with respect to the ground state were calculated from f748pn effective interactions with experimental excitation energies of <sup>45</sup>Ti

Exp .Res[10]		F748pn	J <sup>π</sup>
J <sup>π</sup>	Energy elevels		
5/2 <sup>-</sup>	0.039	0.000	5/2 <sub>1</sub> <sup>-</sup>
7/2 <sup>-</sup>	0.000	0.210	7/2 <sub>1</sub> <sup>-</sup>
9/2 <sup>-</sup>	1.353	1.411	9/2 <sub>1</sub> <sup>-</sup>
11/2 <sup>-</sup>	1.466	1.490	11/2 <sub>1</sub> <sup>-</sup>
(1/2 <sup>-</sup> TO 7/2 <sup>-</sup> )	1.799	1.729	3/2 <sub>1</sub> <sup>-</sup>
3/2 TO 11/2	2.432	2.477	3/2 <sub>2</sub> <sup>-</sup>
5/2 <sup>-</sup> , 7/2 <sup>-</sup>	2.500	2.544	7/2 <sub>2</sub> <sup>-</sup>
-----	-----	2.927	5/2 <sub>2</sub> <sup>-</sup>
15/2 <sup>-</sup>	3.015	2.990	15/2 <sub>1</sub> <sup>-</sup>
-----	3.200	3.299	1/2 <sub>1</sub> <sup>-</sup>
-----	3.200	3.310	13/2 <sub>1</sub> <sup>-</sup>
-----	3.200	3.334	7/2 <sub>3</sub> <sup>-</sup>
-----	3.200	3.341	9/2 <sub>2</sub> <sup>-</sup>
17/2 <sup>-</sup>	3.601	3.383	17/2 <sub>1</sub> <sup>-</sup>
-----	-----	3.765	9/2 <sub>3</sub> <sup>-</sup>
(11/2 TO 15/2)	3.937	3.856	11/2 <sub>2</sub> <sup>-</sup>
-----	-----	3.892	19/2 <sub>1</sub> <sup>-</sup>
-----	-----	3.911	5/2 <sub>3</sub> <sup>-</sup>
(11/2 TO 15/2)	3.937	4.000	13/2 <sub>2</sub> <sup>-</sup>
-----	-----	4.045	7/2 <sub>4</sub> <sup>-</sup>
(7/2)-	4.723	4.319	7/2 <sub>5</sub> <sup>-</sup>
-----	-----	4.351	11/2 <sub>3</sub> <sup>-</sup>
-----	-----	4.650	15/2 <sub>2</sub> <sup>-</sup>
-----	-----	4.818	9/2 <sub>4</sub> <sup>-</sup>
-----	-----	4.844	11/2 <sub>4</sub> <sup>-</sup>
(17/2+)	5.239	5.114	17/2 <sub>2</sub> <sup>-</sup>
-----	5.540	5.120	5/2 <sub>4</sub> <sup>-</sup>
-----	5.540	5.163	15/2 <sub>3</sub> <sup>-</sup>
-----	5.540	5.167	11/2 <sub>5</sub> <sup>-</sup>
-----	5.540	5.167	9/2 <sub>5</sub> <sup>-</sup>
-----	5.540	5.190	13/2 <sub>3</sub> <sup>-</sup>

-----	5.540	5.234	$3/2_3^-$
-----	5.540	5.365	$19/2_2^-$
-----	5.540	5.465	$7/2_6^-$
-----	5.540	5.494	$5/2_5^-$
-----	5.540	5.503	$13/2_4^-$
-----	5.540	5.580	$11/2_6^-$
-----	5.540	5.626	$9/2_6^-$
-----	5.540	5.708	$1/2_2^-$
-----	5.540	5.766	$21/2_1^-$
-----	5.540	5.770	$3/2_4^-$
-----	5.540	5.784	$7/2_7^-$
$23/2^-$	6.163	5.814	$23/2_1^-$
-----	-----	5.974	$5/2_6^-$
-----	-----	6.064	$1/2_3^-$
-----	-----	6.071	$15/2_4^-$
-----	-----	6.141	$3/2_5^-$
-----	-----	6.143	$13/2_5^-$
-----	-----	6.160	$17/2_3^-$
-----	-----	6.179	$9/2_7^-$
-----	-----	6.195	$11/2_7^-$
-----	-----	6.211	$15/2_5^-$
-----	-----	6.216	$7/2_8^-$
-----	-----	6.295	$17/2_4^-$
-----	-----	6.344	$13/2_6^-$
-----	-----	6.371	$9/2_8^-$
-----	-----	6.405	$5/2_7^-$
-----	-----	6.457	$3/2_6^-$
-----	-----	6.536	$13/2_7^-$
-----	-----	6.657	$11/2_8^-$
-----	-----	6.668	$15/2_6^-$
-----	-----	6.677	$7/2_9^-$
-----	-----	6.689	$13/2_8^-$
-----	-----	6.725	$9/2_8^-$
-----	-----	6.789	$19/2_3^-$
-----	-----	6.843	$7/2_{10}^-$
-----	-----	6.857	$15/2_7^-$
-----	-----	6.937	$11/2_9^-$
$27/2^-$	7.143	6.881	$27/2_1^-$
-----	-----	6.956	$5/2_8^-$
-----	-----	6.977	$21/2_2^-$
-----	-----	7.047	$11/2_{10}^-$
-----	-----	7.126	$17/2_5^-$
-----	-----	7.187	$9/2_9^-$
-----	-----	7.403	$17/2_6^-$
-----	-----	7.404	$19/2_4^-$
-----	-----	7.432	$15/2_8^-$
-----	-----	7.441	$13/2_9^-$

-----	-----	7.505	13/2 <sub>10</sub> <sup>-</sup>
-----	7.830	7.543	23/2 <sub>2</sub> <sup>-</sup>
-----	7.830	7.588	15/2 <sub>9</sub> <sup>-</sup>
-----	7.830	7.703	3/2 <sub>7</sub> <sup>-</sup>
-----	7.830	7.709	1/2 <sub>4</sub> <sup>-</sup>
-----	7.830	7.714	25/2 <sub>1</sub> <sup>-</sup>
-----	7.830	7.732	17/2 <sub>7</sub> <sup>-</sup>
-----	7.830	7.841	5/2 <sub>9</sub> <sup>-</sup>
-----	7.830	7.879	19/2 <sub>5</sub> <sup>-</sup>
-----	7.830	7.908	15/2 <sub>10</sub> <sup>-</sup>
-----	7.830	7.939	3/2 <sub>8</sub> <sup>-</sup>
-----	-----	8.354	5/2 <sub>10</sub> <sup>-</sup>
-----	-----	8.574	21/2 <sub>3</sub> <sup>-</sup>
-----	-----	8.840	17/2 <sub>8</sub> <sup>-</sup>
-----	-----	9.009	19/2 <sub>6</sub> <sup>-</sup>
-----	-----	9.013	1/2 <sub>5</sub> <sup>-</sup>
-----	-----	9.069	3/2 <sub>9</sub> <sup>-</sup>
-----	9.643	9.650	19/2 <sub>7</sub> <sup>-</sup>
-----	9.643	9.693	21/2 <sub>4</sub> <sup>-</sup>
-----	9.643	9.706	23/2 <sub>3</sub> <sup>-</sup>
-----	9.643	9.764	3/2 <sub>10</sub> <sup>-</sup>
-----	-----	10.657	17/2 <sub>9</sub> <sup>-</sup>

For <sup>46</sup>Ti nucleus the energy levels calculations from f748pn effective interaction is shown in table 6. The effective interactions give results reasonably consistent with experimental data. The total angular momentum and parity are (0<sub>1</sub><sup>+</sup>, 2<sub>1</sub><sup>+</sup>, 4<sub>1</sub><sup>+</sup>, 2<sub>2</sub><sup>+</sup>, 6<sub>1</sub><sup>+</sup>, 1<sub>1</sub><sup>+</sup>, 2<sub>5</sub><sup>+</sup>, 4<sub>6</sub><sup>+</sup>, 8<sub>3</sub><sup>+</sup>, 4<sub>9</sub><sup>+</sup>, 11<sub>2</sub><sup>+</sup>, 1<sub>5</sub><sup>+</sup>, 12<sub>3</sub><sup>+</sup>, 1<sub>7</sub><sup>+</sup>) respectively, confirmation of which is ( 4<sup>+</sup>)

**Table 6** shows a comparison of the energy levels values with respect to the ground state were calculated from f748pn effective interactions with experimental excitation energies of <sup>46</sup>Ti

Exp .Res[11]		F748pn	J <sup>π</sup>
J <sup>π</sup>	Energy elevels		
0 <sup>+</sup>	0.000	0.000	0 <sub>1</sub> <sup>+</sup>
2 <sup>+</sup>	0.889	0.945	2 <sub>1</sub> <sup>+</sup>
4 <sup>+</sup>	2.009	1.890	4 <sub>1</sub> <sup>+</sup>
2 <sup>+</sup>	2.961	2.541	2 <sub>2</sub> <sup>+</sup>
-----	3.213	2.815	4 <sub>2</sub> <sup>+</sup>
6 <sup>+</sup>	3.298	2.855	6 <sub>1</sub> <sup>+</sup>
-----	3.338	3.306	2 <sub>3</sub> <sup>+</sup>
-----	3.553	3.453	3 <sub>1</sub> <sup>+</sup>
-----	3.553	3.458	4 <sub>3</sub> <sup>+</sup>
-----	3.579	3.540	5 <sub>1</sub> <sup>+</sup>
1 <sup>+</sup>	3.731	3.662	1 <sub>1</sub> <sup>+</sup>
+	3.771	3.703	6 <sub>2</sub> <sup>+</sup>
-----	4.322	4.249	6 <sub>3</sub> <sup>+</sup>
-----	4.322	4.276	8 <sub>1</sub> <sup>+</sup>

-----	4.617	4.642	5 <sub>2</sub> <sup>+</sup>
-----	4.617	4.695	4 <sub>4</sub> <sup>+</sup>
+	4.845	4.851	0 <sub>2</sub> <sup>+</sup>
-----	5.000	4.907	3 <sub>2</sub> <sup>+</sup>
(4 <sup>+</sup> )	5.079	4.957	4 <sub>5</sub> <sup>+</sup>
+	5.094	4.961	3 <sub>3</sub> <sup>+</sup>
+	5.094	5.028	7 <sub>1</sub> <sup>+</sup>
+	5.094	5.068	2 <sub>4</sub> <sup>+</sup>
-----	5.154	5.161	7 <sub>2</sub> <sup>+</sup>
+	5.180	5.272	8 <sub>2</sub> <sup>+</sup>
2 <sup>+</sup>	5.363	5.364	2 <sub>5</sub> <sup>+</sup>
-----	-----	5.420	5 <sub>3</sub> <sup>+</sup>
-----	-----	5.424	6 <sub>4</sub> <sup>+</sup>
-----	-----	5.504	10 <sub>1</sub> <sup>+</sup>
4 <sup>+</sup>	5.794	5.513	4 <sub>6</sub> <sup>+</sup>
+	5.811	5.696	7 <sub>3</sub> <sup>+</sup>
+	5.840	5.839	6 <sub>5</sub> <sup>+</sup>
+	5.903	5.896	5 <sub>4</sub> <sup>+</sup>
-----	6.025	6.033	9 <sub>1</sub> <sup>+</sup>
-----	6.025	6.064	6 <sub>6</sub> <sup>+</sup>
-----	6.025	6.086	0 <sub>3</sub> <sup>+</sup>
-----	-----	6.097	1 <sub>2</sub> <sup>+</sup>
8 <sup>+</sup>	6.200	6.109	8 <sub>3</sub> <sup>+</sup>
-----	6.251	6.225	3 <sub>4</sub> <sup>+</sup>
-----	6.305	6.302	2 <sub>6</sub> <sup>+</sup>
-----	6.305	6.354	10 <sub>2</sub> <sup>+</sup>
+	6.424	6.378	5 <sub>5</sub> <sup>+</sup>
-----	6.513	6.493	6 <sub>7</sub> <sup>+</sup>
-----	6.513	6.519	4 <sub>7</sub> <sup>+</sup>
+	6.550	6.562	8 <sub>4</sub> <sup>+</sup>
-----	6.574	6.648	4 <sub>8</sub> <sup>+</sup>
+	6.616	6.653	11 <sub>1</sub> <sup>+</sup>
4 <sup>+</sup>	6.685	6.696	4 <sub>9</sub> <sup>+</sup>
-----	-----	6.700	3 <sub>5</sub> <sup>+</sup>
-----	-----	6.716	7 <sub>4</sub> <sup>+</sup>
-----	6.794	6.761	9 <sub>2</sub> <sup>+</sup>
-----	6.794	6.795	2 <sub>7</sub> <sup>+</sup>
+	6.851	6.805	8 <sub>5</sub> <sup>+</sup>
+	6.851	6.850	6 <sub>8</sub> <sup>+</sup>
+	6.974	6.890	9 <sub>3</sub> <sup>+</sup>
+	6.974	6.893	5 <sub>6</sub> <sup>+</sup>
+	7.041	7.035	3 <sub>6</sub> <sup>+</sup>
+	7.041	7.048	12 <sub>1</sub> <sup>+</sup>
+	7.101	7.097	7 <sub>5</sub> <sup>+</sup>
+	7.147	7.156	6 <sub>9</sub> <sup>+</sup>
-----	7.172	7.191	10 <sub>3</sub> <sup>+</sup>
-----	7.238	7.259	5 <sub>7</sub> <sup>+</sup>



	+	7.288	7.268	4 <sub>10</sub> <sup>+</sup>
	+	7.288	7.299	2 <sub>8</sub> <sup>+</sup>
	+	7.350	7.377	3 <sub>7</sub> <sup>+</sup>
	+	7.350	7.379	8 <sub>6</sub> <sup>+</sup>
	+	7.558	7.561	5 <sub>8</sub> <sup>+</sup>
	+	7.558	7.562	8 <sub>7</sub> <sup>+</sup>
	-----	7.660	7.688	6 <sub>10</sub> <sup>+</sup>
	+	7.710	7.708	7 <sub>6</sub> <sup>+</sup>
	+	7.788	7.798	2 <sub>9</sub> <sup>+</sup>
	+	7.849	7.801	10 <sub>4</sub> <sup>+</sup>
	+	7.849	7.839	5 <sub>9</sub> <sup>+</sup>
	-----	7.874	7.884	3 <sub>8</sub> <sup>+</sup>
	-----	7.937	7.953	5 <sub>10</sub> <sup>+</sup>
	-----	8.013	8.007	8 <sub>8</sub> <sup>+</sup>
	-----	8.013	8.011	9 <sub>3</sub> <sup>+</sup>
	-----	8.134	8.118	7 <sub>7</sub> <sup>+</sup>
	+	8.182	8.191	3 <sub>9</sub> <sup>+</sup>
	+	8.182	8.204	0 <sub>4</sub> <sup>+</sup>
	+	8.230	8.213	1 <sub>3</sub> <sup>+</sup>
	+	8.230	8.245	2 <sub>10</sub> <sup>+</sup>
	10,11,12+	8.283	8.383	11 <sub>2</sub> <sup>+</sup>
	+	8.384	8.400	9 <sub>4</sub> <sup>+</sup>
	+	8.467	8.467	8 <sub>9</sub> <sup>+</sup>
	+	8.574	8.557	7 <sub>8</sub> <sup>+</sup>
	+	8.621	8.609	3 <sub>10</sub> <sup>+</sup>
	+	8.701	8.744	1 <sub>4</sub> <sup>+</sup>
	+	8.761	8.759	7 <sub>9</sub> <sup>+</sup>
	+	8.808	8.776	10 <sub>4</sub> <sup>+</sup>
	+	8.860	8.865	7 <sub>10</sub> <sup>+</sup>
	+	8.860	8.879	8 <sub>10</sub> <sup>+</sup>
	+	8.984	8.984	10 <sub>5</sub> <sup>+</sup>
	+	9.070	9.085	12 <sub>2</sub> <sup>+</sup>
	+	9.253	9.202	14 <sub>1</sub> <sup>+</sup>
	+	9.253	9.246	9 <sub>5</sub> <sup>+</sup>
	+	9.345	9.332	13 <sub>1</sub> <sup>+</sup>
	1 <sup>+</sup>	9.420	9.441	1 <sub>5</sub> <sup>+</sup>
	+	9.474	9.467	0 <sub>5</sub> <sup>+</sup>
	+	9.572	9.536	11 <sub>2</sub> <sup>+</sup>
	12 <sup>+</sup> , 14 <sup>+</sup>	10.041	9.927	12 <sub>3</sub> <sup>+</sup>
	-----	10.212	10.141	9 <sub>6</sub> <sup>+</sup>
	-----	10.256	10.271	1 <sub>6</sub> <sup>+</sup>
	-----	10.321	10.300	10 <sub>6</sub> <sup>+</sup>
	+	10.523	10.525	9 <sub>7</sub> <sup>+</sup>
	-----	11.426	11.456	9 <sub>8</sub> <sup>+</sup>
	1 <sup>+</sup>	11.450	11.484	1 <sub>7</sub> <sup>+</sup>
	-----	-----	11.577	10 <sub>7</sub> <sup>+</sup>
	-----	-----	11.981	11 <sub>3</sub> <sup>+</sup>

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12.974

12.835

0<sub>6</sub><sup>+</sup>

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## References:

- [1] B. A. Brown, C. R. Bronk and P. E. Hodgson, "Systematics of nuclear RMS charge radii." , *Nuclear Physics.*, **10(12)** , 1683 , 1984.
- [2] B. A. Brown, S. E. Massen and J. I. Escudero, "PE llogson, G. Madurga and J. Vinas." , *J. Phys.*, **G 9**, 423,1983.
- [3] S. Cohen and D. Kurath, "Effective interactions for the 1p shell." , *Nuclear Physics.*, **73(1)** , 1-24 , 1965.
- [4] B. Alex. Brown and B. H. Wildenthal , "Status of the nuclear shell model." , *Annual Review of Nuclear and Particle Science.*, **38(1)** , 29-66 ,1988.
- [5] A. B. Brown and R. Sherr "Charge-dependent two body interactions deduced from displacement energies in the 1f<sub>7/2</sub>shell", *Nucl . Phy.*, **A 322(1)**, 61-91 ,1979.
- [6] A. Gade, D. Bazin, C. A. Bertulani, B. A. Brwon, C. A. Church, D. C. Dinca, J. Enders, T. P. Glasmacher, G. Hansen, Z. Hu, K. W. Kemper, W. F. Muller, H. Olliver, B. C. Perry, L. A. Riley, B. T. Roeder, B. M. Sherrill, J. R. Terry, J. A. Tostevin and K. L. Yurkewics, *Phys. Rev.*, **C71** , 051301(R) ,2005.
- [7] S. J. Freeman, R. V. F. Janssens, B. A. Brown, M. P. Carpenter, S. M. Fischer, N. J. Hammond, M. Honma, C. J. Lister, T. L. Khoo, G. Mkherjee, D. Seweryaniak, J. F. Smith, B. J. Varley, M. whitehead and S. Zhu., *Phys. Rev.*, **764** , 142-147 ,2005.
- [8] L. Rydström, J. Blomqvist, R. J. Liotta and C. Pomar, Structure of proton-deficient nuclei near 208Pb. *Nuclear Physics A* , **512(2)** , 217-240 ,1990.
- [9] B. A. Brown, A. Etchegoyen , N. S. Godwin, W. D. M. Rae, W. A. Richter, W. E. Ormand , E. K. Warburton, J. S. Winfield, L. Zhao and C. H. Zimmerman , *MSU-NSCL report number.*,1289 ,2004 .
- [10] T.W.Burrows, " Nuclear Data Sheets for A = 45 " , *Nuclear Data Sheets.*, **109(1)**,171-296 ,2008.
- [11]S.C.Wu, " Nuclear Data Sheets for A = 46 " , *Nuclear Data Sheets.*, **91(1)**, 1-116 ,2000.