

# NONHOLONOMIC FRAMES FOR FINSLER SPACE WITH SPECIAL $(\alpha, \beta)$ -METRICS

**NARASIMHAMURTHY S.K., MALLIKARJUN Y. KUMBAR., KAVYASHREE A.R.**

Department of P.G. Studies and Research in Mathematics,

Kuvempu University, Shankaraghatta - 577451,

Shimoga, Karnataka, INDIA

E-mail: nmurthysk@hotmail.com and [mallikarjunykumbar@gmail.com](mailto:mallikarjunykumbar@gmail.com)

**Abstract:** The main aim of this paper is to, first we determine the two Finsler deformations to the special Finsler  $(\alpha, \beta)$ -metrics. Consequently, we obtain the nonholonomic frame for the special  $(\alpha, \beta)$ -metrics, such as  $L = \alpha + \beta + \frac{\beta^2}{\alpha}$  (1<sup>st</sup> Approximate Matsumoto metric) and  $L = \frac{\beta^2}{\beta - \alpha}$  (Infinite series  $(\alpha, \beta)$ -metric).

**Key Words:** *Finsler Space,  $(\alpha, \beta)$ -metrics, GL-metric, Nonholonomic Finsler frame.*

## 1. INTRODUCTION

In 1982, P.R. Holland ([1][2]), studies a unified formalism that uses a nonholonomic frame on space-time arising from consideration of a charged particle moving in an external electromagnetic field. In fact, R.S. Ingarden [3] was first to point out that the Lorentz force law can be written in this case as geodesic equation on a Finsler space called Randers space. The author Beil R.G. ([5][6]), have studied a gauge transformation viewed as a nonholonomic frame on the tangent bundle of a four dimensional base manifold. The geometry that follows from these considerations gives a unified approach to gravitation and gauge symmetries. The above authors used the common Finsler idea to study the existence of a nonholonomic frame on the vertical subbundle  $V TM$  of the tangent bundle of a base manifold  $M$ .

In this paper, the fundamental tensor field might be taught as the result of two Finsler deformation. Then we can determine a corresponding frame for each of these two Finsler deformations. Consequently, a nonholonomic frame for a Finsler space with special  $(\alpha, \beta)$ -metrics such as Ist Approximate Matsumoto metric and Infinite series  $(\alpha, \beta)$ -metrics will appear as a product of two Finsler frames formerly determined. This is an extension work of Ioan Bucataru and Radu Miron [10].

Consider  $a_{ij}(x)$ , the components of a Riemannian metric on the base manifold  $M$ ,  $a(x, y) > 0$  and  $b(x, y) \geq 0$  two functions on  $TM$  and  $B(x, y) = B_i(x, y)dx^i$  a vertical 1-form on  $TM$ . Then

$$g_{ij}(x, y) = a(x, y)a_{ij}(x) + b(x, y)B_i(x, y)B_j(x, y) \quad (1.1)$$

is a generalized Lagrange metric, called the *Beil metric*. We say also that the metric tensor  $g_{ij}$  is a *Beil deformation* of the Riemannian metric  $a_{ij}$ . It has been studied and applied by R.Miron and R.K. Tavakol in General Relativity for  $a(x, y) = \exp(2\sigma(x, y))$  and  $b = 0$ . The case  $a(x, y) = 1$  with various choices of  $b$  and  $B_i$  was introduced and studied by R.G. Beil for constructing a new unified field theory [6].

## 2. PRELIMINARIES

An important class of Finsler spaces is the class of Finsler spaces with  $(\alpha, \beta)$ -metrics [11]. The first Finsler spaces with  $(\alpha, \beta)$ -metrics were introduced by the physicist G.Randers in 1940, are called Randers spaces [4]. Recently, R.G. Beil suggested to consider a more general case, the class of Lagrange spaces with  $(\alpha, \beta)$ -metric, which was discussed in [12]. A unified formalism which uses a nonholonomic frame on space time, a sort of plastic deformation, arising from consideration of a charged particle moving in an external electromagnetic field in the background space time viewed as a strained mechanism studied by P.R.Holland [1][2]. If we do not ask for the function  $L$  to be homogeneous of order two with respect to the  $(\alpha, \beta)$  variables, then we have a *Lagrange space with  $(\alpha, \beta)$ -metric*. Next we look for some different Finsler space with  $(\alpha, \beta)$ -metrics.

**Definition 2.1:** A Finsler space  $F^n = (M, F(x, y))$  is called with  $(\alpha, \beta)$ -metric if there exists a 2-homogeneous function  $L$  of two variables such that the Finsler metric  $F: TM \rightarrow R$  is given by,

$$F^2(x, y) = L(\alpha(x, y), \beta(x, y)),$$

Where  $\alpha^2(x, y) = a_{ij}(x)y^i y^j$ ,  $\alpha$  is a Riemannian metric on  $M$ , and (2.1)

$\beta(x, y) = b_i(x)y^i$  is a 1-form on  $M$ .

Consider  $g_{ij} = \frac{1}{2} \frac{\partial^2 F^2}{\partial y^i \partial y^j}$  the fundamental tensor of the Randers space  $(M, F)$ . Taking into account the homogeneity of  $\alpha$  and  $F$  we have the following formulae:

$$p^i = \frac{1}{\alpha} y^i = a^{ij} \frac{\partial \alpha}{\partial y^j} ; \quad p_i = a_{ij} p^j = \frac{\partial \alpha}{\partial y^i} ;$$

$$l^i = \frac{1}{L} y^i = g^{ij} \frac{\partial L}{\partial y^j} ; \quad l_i = g_{ij} l^j = \frac{\partial L}{\partial y^i} = P_i + b_i; \quad (2.2)$$

$$l^i = \frac{1}{L} p^i ; \quad l^i l_i = P^i p_i = 1; \quad l^i p_i = \frac{\alpha}{L}; \quad p^i l_i = \frac{L}{\alpha};$$

$$b_i P^i = \frac{\beta}{\alpha} ; \quad b_i l^i = \frac{\beta}{L} .$$

With respect to these notations, the metric tensors  $(a_{ij})$  and  $(g_{ij})$  are related by [13],

$$g_{ij}(x, y) = \frac{L}{\alpha} a_{ij} + b_i P_j + P_i b_j + b_i b_j - \frac{\beta}{\alpha} p_i p_j = \frac{L}{\alpha} (a_{ij} - p_i p_j) + l_i l_j. \quad (2.3)$$

**Theorem 2.1[10]:** For a Finsler space  $(M, F)$  consider the matrix with the entries:

$$Y_j^i = \sqrt{\frac{\alpha}{L}} \left( \delta_j^i - l^i l_j + \sqrt{\frac{\alpha}{L}} p^i p_j \right) \quad (2.4)$$

defined on  $TM$ . Then  $Y_j = Y_j^i \left( \frac{\partial}{\partial y^i} \right)$ ,  $j \in \{1, 2, \dots, n\}$  is an nonholonomic frame.

**Theorem 2.2 [7]:** With respect to frame the holonomic components of the Finsler metric tensor  $(a_{\alpha\beta})$  is the Randers metric  $(g_{ij})$ , i.e.,

$$g_{ij} = Y_i^\alpha Y_j^\beta a_{\alpha\beta}. \quad (2.5)$$

Throughout this section we shall rise and lower indices only with the Riemannian metric  $a_{ij}(x)$  that is  $y_i = a_{ij} y^j$ ,  $b^i = a^{ij} b_j$ , and so on. For a Finsler space with  $(\alpha, \beta)$ -metric  $F^2(x, y) = L(\alpha(x, y), \beta(x, y))$  we have the Finsler invariants [13].

$$\rho_1 = \frac{1}{2\alpha} \frac{\partial L}{\partial \alpha}; \quad \rho_0 = \frac{1}{2} \frac{\partial^2 L}{\partial \beta^2}; \quad \rho_{-1} = \frac{1}{2\alpha} \frac{\partial^2 L}{\partial \alpha \partial \beta}; \quad \rho_{-2} = \frac{1}{2\alpha^2} \left( \frac{\partial^2 L}{\partial \alpha^2} - \frac{1}{\alpha} \frac{\partial L}{\partial \alpha} \right) \quad (2.6)$$

where subscripts 1, 0, -1, -2 gives us the degree of homogeneity of these invariants.

For a Finsler space with  $(\alpha, \beta)$ -metric we have,

$$\rho_{-1} \beta + \rho_{-2} \alpha^2 = 0 \quad (2.7)$$

with respect to the notations we have that the metric tensor  $g_{ij}$  of a Finsler space with  $(\alpha, \beta)$ -metric is given by [13].

$$g_{ij}(x, y) = \rho a_{ij}(x) + \rho_0 b_i(x) + \rho_{-1} (b_i(x) y_j + b_j(x) y_i) + \rho_{-2} y_i y_j. \quad (2.8)$$

From (2.8) we can see that  $g_{ij}$  is the result of two Finsler deformations:

$$\left. \begin{aligned} \text{i)} \quad a_{ij} \rightarrow h_{ij} &= \rho a_{ij} + \frac{1}{\rho_{-2}} (\rho_{-1} b_i + \rho_{-2} y_i) (\rho_{-1} b_j + \rho_{-2} y_j) \\ \text{ii)} \quad h_{ij} \rightarrow g_{ij} &= h_{ij} + \frac{1}{\rho_{-2}} (\rho_0 \rho_{-2} - \rho_{-1}^2) b_i b_j. \end{aligned} \right\} \quad (2.9)$$

The nonholonomic Finsler frame that corresponding to the 1<sup>st</sup> deformation (2.9) is according to the theorem (7.9.1) in [10], given by,

$$X_j^i = \sqrt{\rho} \delta_j^i - \frac{1}{\beta^2} \left( \sqrt{\rho} + \sqrt{\rho + \frac{\beta^2}{\rho_{-2}}} \right) (\rho_{-1} b^i + \rho_{-2} y^i) (\rho_{-1} b^i + \rho_{-2} y^j) \quad (2.10)$$

where  $B^2 = a_{ij} (\rho_{-1} b^i + \rho_{-2} y^i) (\rho_{-1} b^i + \rho_{-2} y^j) = \rho_{-1}^2 b^2 + \beta \rho_{-1} \rho_{-2}$

This metric tensor  $a_{ij}$  and  $h_{ij}$  are related by,

$$h_{ij} = X_i^k X_j^l a_{kl}. \quad (2.11)$$

again the frame that corresponds to the  $II^{nd}$  deformation (2.9) given by,

$$Y_j^i = \delta_j^i - \frac{1}{c^2} \left( 1 \pm \sqrt{1 + \left( \frac{\rho_{-2} c^2}{\rho_0 \rho_{-2} - \rho_{-1}^2} \right)} \right) b^i b_j, \quad (2.12)$$

where  $C^2 = h_{ij} b^i b^j = \rho b^2 + \frac{1}{\rho_{-2}} (\rho_{-1} b^2 + \rho_{-2} \beta)^2$ .

The metric tensor  $h_{ij}$  and  $g_{ij}$  are related by the formula;

$$g_{mn} = Y_m^i Y_n^j h_{ij}. \quad (2.13)$$

**Theorem 2.3:** [10] Let  $F^2(x, y) = L(\alpha(x, y), \beta(x, y))$  be the metric function of a Finsler space with  $(\alpha, \beta)$ -metric for which the condition (2.7) is true. Then

$$V_j^i = X_k^i Y_j^k$$

is a nonholonomic Finsler frame with  $X_k^i$  and  $Y_j^k$  are given by (2.10) and (2.12) respectively.

### 3. NONHOLONOMIC FRAMES FOR FINSLER SPACE WITH SPECIAL $(\alpha, \beta)$ -METRICS

In this section we consider two Finsler space with special  $(\alpha, \beta)$ -metrics, such as Ist Approximate Matusmoto metric and Infinite series  $(\alpha, \beta)$ -metric, then we construct nonholonomic Finsler frames.

### 3.1. NONHOLONOMIC FRAME FOR $L = \left(\alpha + \beta + \frac{\beta^2}{\alpha}\right)^2$ :

In the first case, for a Finsler space with the fundamental function  $L = \left(\alpha + \beta + \frac{\beta^2}{\alpha}\right)^2$ , the Finsler invariants (2.6) are given by:

$$\left. \begin{aligned} \rho_1 &= \frac{(\alpha^2 + \alpha\beta + \beta^2)(\alpha^2 - \beta^2)}{\alpha^4}, & \rho_0 &= \frac{3(\alpha^2 + 2\alpha\beta + 2\beta^2)}{\alpha^2}, \\ \rho_{-1} &= \frac{(\alpha^3 - 3\alpha\beta^2 - 4\beta^3)}{\alpha^4}, & \rho_{-2} &= -\frac{\beta(\alpha^3 - 3\alpha\beta^2 - 4\beta^3)}{\alpha^6}, \\ B^2 &= \frac{(\alpha^3 - 3\alpha\beta^2 - 4\beta^3)^2(b^2\alpha^2 - \beta^2)}{\alpha^{10}}. \end{aligned} \right\} \quad (3.1)$$

Using (3.1) in (2.10) we have,

$$\begin{aligned} X_j^i &= \sqrt{\frac{(\alpha^2 + \alpha\beta + \beta^2)(\alpha^2 - \beta^2)}{\alpha^4}} \delta_j^i \\ &\quad - \frac{\alpha^4}{b^2\alpha^2 - \beta^2} \left\{ \sqrt{(\alpha^2 + \alpha\beta + \beta^2)(\alpha^2 - \beta^2)} \pm \sqrt{(\alpha^2 + \alpha\beta + \beta^2)(\alpha^2 - \beta^2) - \frac{\alpha^3 - 3\alpha\beta^2 - 4\beta^3}{\beta}} \right\} \\ &\quad \left( b^i - \frac{\beta y^i}{\alpha^2} \right) \left( b_j - \frac{\beta y_j}{\alpha^2} \right); \end{aligned} \quad (3.2)$$

Again using (3.1) in (2.12) we have,

$$Y_j^i = \delta_j^i - \frac{1}{c^2} \left( 1 \pm \sqrt{1 + \frac{\alpha^2 \beta C^2}{\alpha^3 + 3\alpha\beta(\alpha + \beta) + 2\beta^3}} \right) b^i b_j; \quad (3.3)$$

where 
$$C^2 = \frac{(\alpha^6 + \alpha\beta + \beta^2)(\alpha^2 - \beta^2)b^2}{\alpha^4} - \frac{(\alpha^3 - 3\alpha\beta^2 - 4\beta^3)(b^2\alpha^2 - \beta^2)^2}{\alpha^6\beta}.$$

**Theorem 3.1:** Consider a Finsler space  $L = \left(\alpha + \beta + \frac{\beta^2}{\alpha}\right)^2$ , for which the condition (2.7) is true. Then

$$V_j^i = X_k^i Y_j^k$$

is a nonholonomic Finsler frame with  $X_k^i$  and  $Y_j^k$  are given by (3.2) and (3.3) respectively.

### 3.2. NONHOLONOMIC FRAME FOR $L = \frac{\beta^4}{(\beta - \alpha)^2}$ :

In the second case, for a Finsler space with the fundamental function  $L = \frac{\beta^4}{(\beta-\alpha)^2}$ , the Finsler invariants (2.6) are given by:

$$\left. \begin{aligned} \rho_1 &= \frac{\beta^4}{\alpha(\beta-\alpha)^3}, \quad \rho_0 = \frac{\beta^2(\beta^2 - 4\alpha\beta + 6\alpha^2)}{(\alpha-\beta)^4}, \\ \rho_{-1} &= \frac{\beta^3(\beta-4\alpha)}{(\alpha-\beta)^4}, \quad \rho_{-2} = \frac{\beta^4(4\alpha-\beta)}{\alpha^3(\alpha-\beta)^4}, \\ B^2 &= \frac{\beta^6(4\alpha-\beta)^2(b^2\alpha^2-\beta)}{\alpha^4(\alpha-\beta)^8}. \end{aligned} \right\} \quad (3.4)$$

Using (3.4) in (2.10) we have,

$$\begin{aligned} X_j^i &= \sqrt{\frac{\beta^4}{\alpha(\beta-\alpha)^3}} \delta_j^i - \frac{\alpha^2}{b^2\alpha^2-\beta^2} \left\{ \sqrt{\frac{\beta^4}{\alpha(\beta-\alpha)^3}} \pm \sqrt{\frac{\beta^2(2\beta^3-5\alpha\beta^2+4b^2\alpha^3-b^2\alpha^2\beta)}{\alpha(\alpha-\beta)^4}} \right\} \\ &\quad \left( b^i - \frac{\beta y^i}{\alpha^2} \right) \left( b_j - \frac{\beta y_j}{\alpha^2} \right); \end{aligned} \quad (3.5)$$

Again using (3.4) in (2.12) we have,

$$Y_j^i = \delta_j^i - \frac{1}{c^2} \left( 1 \pm \sqrt{1 + \frac{(\alpha-\beta)^3 C^2}{\alpha^4}} \right) b^i b_j; \quad (3.6)$$

where 
$$C^2 = \frac{\alpha^2(\alpha-2\beta)b^2}{(\alpha-\beta)^3} - \frac{(\alpha-4\beta)(b^2\alpha^2-\beta^2)^2}{\beta(\alpha-\beta)^4}.$$

**Theorem 3.2:** Consider a Finsler space  $L = \frac{\beta^4}{(\beta-\alpha)^2}$ , for which the condition (2.7) is true. Then

$$V_j^i = X_k^i Y_j^k$$

is a nonholonomic Finsler frame with  $X_k^i$  and  $Y_j^k$  are given by (3.5) and (3.6) respectively.

#### 4. CONCLUSION

Nonholonomic frame relates a semi-Riemannian metric (the Minkowski or the Lorentz metric) with an induced Finsler metric. Antonelli P.L., Bucataru I. ([7][8]), has been determined such a nonholonomic frame for two important classes of Finsler spaces that are dual in the sense of Randers and Kropina spaces [9]. As Randers and Kropina spaces are members of a bigger class of Finsler spaces, namely the Finsler spaces with  $(\alpha, \beta)$ -metric, it appears a natural question: *Does how many Finsler space with  $(\alpha, \beta)$ -metrics have such a nonholonomic frame? The answer is yes, there are many Finsler space with  $(\alpha, \beta)$ -metrics.*

In this work, we consider the two special Finsler metrics and we determine the nonholonomic Finsler frames. Each of the frames we found here induces a Finsler connection on TM with torsion and no curvature. But, in Finsler geometry, there are many  $(\alpha, \beta)$ -metrics, in future work we can determine the frames for them also.

## 5. REFERENCES

- 1) Holland. P.R.: *Electromagnetism, Particles and Anholonomy*. Physics Letters, 91 (6), 275-278 (1982).
- 2) Holland, P.R.: *Anholonomic deformations in the ether: a significance for the electrodynamic potentials*. In: Hiley, B.J. Peat, F.D. (eds.), *Quantum Implications*. Routledge and Kegan Paul, London and New York, 295-311 (1987).
- 3) Ingarden, R.S.: *On Physical interpretations of Finsler and Kawaguchi spaces*. Tensor N.S., 46, 354-360 (1987).
- 4) Randers, G.: *On asymmetric metric in the four space of general relativity*. Phys. Rev., 59, 195-199 (1941).
- 5) Beil, R.G.: *Comparison of unified field theories*. Tensor N.S., 56, 175-183 (1995).
- 6) Beil, R.G.: *Equations of Motion from Finsler Geometric Methods*. In: Antonelli, P.L. (ed), *Finslerian Geometries. A meeting of minds*. Kluwer Academic Publisher, FTPH, no. 109, 95-111 (2000).
- 7) Antonelli, P.L., Bucataru, I.: *On Holland's frame for Randers space and its applications in physics*. In: Kozma, L. (ed), *Steps in Differential Geometry. Proceedings of the Colloquium on Differential Geometry*, Debrecen, Hungary, July 25-30, 2000. Debrecen: Univ. Debrecen, Institute of Mathematics and Informatics, 39-54 (2001).
- 8) Antonelli, P.L., Bucataru, I.: *Finsler connections in anholonomic geometry of a Kropina space*. Nonlinear Studies, 8 (1), 171-184 (2001).
- 9) Hrimiuc, D., Shimada, H.: *On the L-duality between Lagrange and Hamilton manifolds*. Nonlinear World, 3, 613-641 (1996).
- 10) Ioan Bucataru, Radu Miron: *Finsler-Lagrange Geometry. Applications to dynamical systems* CEEEX ET 3174/2005-2007 & CEEEX M III 12595/2007 (2007).
- 11) Matsumoto, M.: *Theory of Finsler spaces with  $(\alpha, \beta)$ -metric*. Rep. Math. Phys., 31, 43-83 (1992).
- 12) Bucataru, I.: *Nonholonomic frames on Finsler geometry*. Balkan Journal of Geometry and its Applications, 7 (1), 13-27 (2002).
- 13) Matsumoto.M: *Foundations of Finsler geometry and special Finsler spaces*, Kaishesha Press, Otsu, Japan, 1986.