

Calculation of Energy Levels and B(E2) for ^{20}Ne Isotope by Using Nuclear Shell Model

Ali k.Hasan and Hadeel H.Abed

* Department of Physics, College of Education for Girls, University of Kufa, Iraq

DOI: 10.29322/IJSRP.9.04.2019.p8861
<http://dx.doi.org/10.29322/IJSRP.9.04.2019.p8861>

Abstract- In this study, the energy levels and the probability of electric transition B(E2) of the ^{20}Ne nucleus were calculated using the OXBASH code within the sd shell and the use of the effective interaction KUOSD. OXBASH is a computing code to perform a nuclear installation calculation based on a shell mode, Energy levels and the probability of electric transition is acceptable agreement with available experimental data.

Index Terms- OXBASH Code, ^{20}Ne nucleus , Energy Levels, B(E2), sd shell

I. INTRODUCTION

One of the criteria to improve investigations of nuclei properties is obtaining the energy levels and nuclear structure of nuclei ,nuclear models have the property to help us to better understanding of nuclear structure which includes main physical properties of nuclei, shell model is one of the most featured and successful nuclear models this model can be compared to the atomic shell model Atomic behaviors and properties can also be described with valence electrons that are outside of a closed shell similarly, nucleon valence (protons or neutrons) in a nucleus is placed outside the closed shell with magic numbers (2,8,20,28,50,82 and 126) play important roles in determining nuclear properties the nucleus with magic numbers is very stable and has completely different properties compared to their neighbors[1]. This model assumes that the nucleons are distributed at separate energy levels according to the Pauli principle, this nucleons move inside a mean field potential produced by the other nucleons[2]. Shell model calculations are performed out within a model space in which the nucleons are limited to a few orbits, when using suitable effective operators, taking into account the effect of the bigger model space ,this model offers a meaningful description of this observable[3].A study of nuclei in the sd shell thus leads to the best understanding between a macroscopic (collective) description and a microscopic description of the nucleus (shell model)[4]. The sd-shell core is considered to be an inert ^{16}O core, on which the valence nucleus is distributed $1d5/2$, $2s1/2$ and $1d3/2$ status[5].

II. THEORY

The calculations of the shell model include single-particle energies (SPE) and two-body interaction matrix elements or called two-body matrix elements (TBME), recently this group was

named (effective interaction) or (model space Hamiltonian). This model space Hamiltonian can be described in two ways: one method is a "realistic" which is created for a particular shell model space of known data on the free nucleon-nucleon force. The second method is "experimental" which is found in the parameters whose values are specified by agreement between measured level energies and shell model eigenvalue [6].For calculations can be written the shell model Hamiltonian as follows[7].

$$H = \sum_i \varepsilon_i n_i + \sum_{ijkl} V_{ijkl} a_i^\dagger a_j^\dagger a_i a_j$$

Where:

ε_i : single particle energy of orbit i, -:

n_i : stands for the number operator of the orbit i

V_{ijkl} : the two body matrix elements for the nucleon-nucleon interaction of orbits(i,j,k,l)

a_i^\dagger : a_j^\dagger , Creation operators

a_i Annihilation operators : a_j ,

It is important to know that there are residual interactions between each of the particles in each case (secondary shells) and that the remaining interactions are observed in the secondary shell if they are partially full, and the angular momentum and parity of the ground state of the nucleus can be determined by the remaining interactions, the nucleic levels of the nucleus have a momentum equal to zero and their positive parity. The nucleons in the secondary shell give double energy in which the values of M_j are inverse and the number of magnetic quantities takes the following values:[8]

$$M_j = j, (j-1), \dots, 1/2, -1/2, \dots, -(j-1), -j$$

The probability of electric transition can be calculated by equation:[9]

$$B(EL: J_i \rightarrow J_f) = \frac{1}{(2J_i + 1)} \sum_{M_i M_f} \left| \langle \psi_{J_f M_f} | T_{LM}^{(E)} | \psi_{J_i M_i} \rangle \right|^2$$

$$\left(\frac{2J_f + 1}{2J_i + 1} \right) \left| \langle \psi_{J_f} | T_L^{(E)} | \psi_{J_i} \rangle \right|^2 =$$

Where $T_L^{(E)}$ is the electric multipole operator , J_i and J_f are the spin of the initial and final states, respectively

III. SHELL MODEL CALCULATION

Calculations of energy levels and B(E2) were made with the code OXBASH[10]. OXBASH is using an m-scheme Slater determinant basis with wave functions and good angular momentum J, isospin T is constructed, this program is the only version that uses the projection technology, the sd model spaces contain (0d5/2, 1s1/2, and 0d3/2) above the N = 8 and Z = 8 closed shells for neutrons and protons [11]. Considering the number of valence nucleons for Neon isotopes, the sd model space is the appropriate area for our calculations, Where it is calculated energy level and B(E2) for 20Ne employing harmonic oscillator potential (HO, b), b<0 for Ne isotope within this space . In this work, we focus our attention on the description of energy levels and B(E2) of sd shell of 20Ne isotopes which have configurations(0d5/2, 1s1/2and 0d3/2) ,the sd model space with KUOSD interaction

3.1 Energy levels of ²⁰Ne nucleus

According to the shell model, the ground state of ²⁰Ne nucleus is a closed ¹⁶O core , with J =0+ and T=0,the excited states consist of the configuration of the nucleons in the sd-shell .in this work we use KUOSD interaction and single-particle energies (SPEs) are {1d5/2=-4.150,2s1/2=-3.280and 1d3/2=0.930} respectively.The table(1) shows a comparison between theoretical results and available experimental results[12] for Neon isotopes by using KUOSD interaction. From the down table KUOSD Hamiltonians agree reasonably well with experimental data .

The ground state was confirmed 0⁺.The agreement is good for the states of with experimental values.

The agreement is good for the states (3₁⁺, 1₁⁺, 2₆⁺, 2₈⁺, 6₃⁺, 2₁⁺, 4₁⁺, 0₂⁺, 6₁⁺, 2₂⁺, 4₂⁺, 2₃⁺, 2₄⁺, 4₃⁺, 1₂⁺, 8₁⁺, 0₃⁺, 4₄⁺, 2₅⁺, 6₂⁺, 4₅⁺, 4₇⁺, 6₃⁺, 4₈⁺, 8₂⁺, 4₉⁺, 0₅⁺, 6₅⁺, 4₁₀⁺, 6₆⁺, 6₇⁺, 6₈⁺, 8₃⁺) as compared with the experimental data. This study also confirmed the angular momentum and parity for energies (11.653,14.455,14,731,14.653,16.010,15.970,17.910)Mev with states (3₂⁺, 2₇⁺, 4₆⁺, 0₄⁺, 2₉⁺, 6₄⁺, 0₆⁺) after comparing them with our theoretical values. The total angular momentum values were confirmed without parity for level 2₁₀⁺.

The values of experimental energies (13.426,14.313,15.436) Mev were confirmed for angular momentum (5₂⁺, 3₅⁺, 3₆⁺) with positive parity.

The agreement is good for angular momentum without parity for state

(5₄⁺, 3₃⁺, 5₃⁺, 7₁⁺, 3₇⁺, 3₈⁺, 3₉⁺, 5₇⁺, 7₃⁺, 5₈⁺, 7₄⁺, 7₅⁺, 3₄⁺, 5₅⁺, 5₆⁺, 7₂⁺, 5₁₀⁺).

The levels of for which the angular momentum and parity are yet unknown at the states

(5₄⁺, 1₅⁺, 1₆⁺, 1₇⁺, 0₇⁺, 1₈⁺, 5₉⁺, 1₉⁺, 0₈⁺, 1₁₀⁺, 0₉⁺, 6₁₀⁺, 0₁₀⁺, 3₁₀⁺)

And through our calculations we noticed that the highest energy value is (24.25)Mev for angular momentum and parity 7₅⁺ of reaction KUOSD, while the highest experimental value is (28.20)Mev.

Table1.Comparison of the experimental excitation energies[12] and excitation energies predictions for ²⁰Ne nucleus by using KUOSD interactions

J ⁺ (OXBASH)	Energy (OXBASH) KUOSD(ME)	Energy(exp) (Mev)	J(exp)
0 ₁	0.0	0.0	0 ⁺
2 ₁	1.455	1.633	2 ⁺
4 ₁	3.861	4.247	4 ⁺
0 ₂	6.81	6.725	0 ⁺
6 ₁	8.188	8.777	6 ⁺
2 ₂	8.434	7.833	2 ⁺
4 ₂	10.181	9.990	4 ⁺
2 ₃	10.374	10.273	2 ⁺
3 ₁	10.692	10.694	4, 3 ⁺
2 ₄	11.097	10.940	2 ⁺
4 ₃	11.188	11.090	4 ⁺
1 ₁	11.781	11.262	1 ⁺
3 ₂	11.904	11.653	(3 ⁺)
8 ₁	12.358	11.951	8 ⁺
5 ₁	12.482	12.713	5 ⁻
0 ₃	12.812	12.436	0 ⁺
4 ₄	12.842	13.048	4 ⁺
2 ₅	13.185	13.095	2 ⁺
5 ₂	13.204	13.426	(5 ⁻)
3 ₃	13.332	13.226	3 ⁻
1 ₂	13.379	13.307	1 ⁺
6 ₂	13.412	13.105	6 ⁺
3 ₄	13.427	13.414	3 ⁻
2 ₆	13.55	13.592	2 ⁺

1 ₃	13.629	13.736	1 ⁺
1 ₄	13.942	14.200	1 ⁺
4 ₅	14.12	14.270	4 ⁺
2 ₇	14.382	14.455	0 ⁺ ,2 ⁺)
5 ₃	14.649	14.816	5 ⁻
4 ₆	14.733	14.731	(4 ⁺)
0 ₄	14.865	14.653	(0 ⁺)
3 ₅	14.931	14.313	(3 ⁻)
7 ₁	14.972	15.366	7 ⁻
2 ₈	15.04	15.047	2 ⁺
4 ₇	15.2	15.330	4 ⁺
6 ₃	15.235	15.159	6 ⁺
3 ₆	15.490	15.436	(3 ⁻)
4 ₈	15.856	16.329	4 ⁺
5 ₄	15.888	-----	-----
3 ₇	15.952	16.136	3 ⁻
2 ₉	15.977	16.010	(2 ⁺)
6 ₄	16.138	15.970	(6 ⁺)
2 ₁₀	16.324	16.437	(0,2,4) ⁺
8 ₂	16.374	16.746	8 ⁺
3 ₈	16.46	16.628	3 ⁻
5 ₅	16.566	16.559	5 ⁻
4 ₉	16.592	16.667	4 ⁺
1 ₅	16.829	-----	-----
3 ₉	17.041	17.284	3 ⁻
0 ₅	17.155	16.732	0 ⁺
6 ₅	17.383	17.541	6 ⁺
0 ₆	17.568	17.910	(0 ⁺)
4 ₁₀	17.624	17.769	4 ⁺
5 ₆	17.895	17.851	5 ⁻
1 ₆	17.941	-----	-----
7 ₂	18.037	18.005	7 ⁻
3 ₁₀	18.038	-----	-----
6 ₆	18.174	18.286	6 ⁺
5 ₇	18.733	18.494	5 ⁻
7 ₃	18.848	18.768	7 ⁻
1 ₇	18.873	-----	-----
0 ₇	18.94	-----	-----
6 ₇	18.988	18.745	6 ⁺
1 ₈	19.332	-----	-----
5 ₈	19.412	19.859	5 ⁻
5 ₉	19.624	-----	-----
6 ₈	19.736	19.655	6 ⁺
1 ₉	20.264	-----	-----
0 ₈	20.45	-----	-----
6 ₉	20.466	20.445	6 ⁺
5 ¹⁰	20.476	20.468	5 ⁻
7 ₄	20.78	20.760	7 ⁻
1 ₁₀	21.023	-----	-----
0 ₉	21.075	-----	-----
6 ₁₀	21.802	-----	-----
0 ₁₀	22.144	-----	-----
8 ₃	22.771	23.40	8 ⁺
7 ₅	24.254	22.30	7 ⁻

3.2 Reduced electric quadrupole transition probability B(E2) Calculation:

The B(E2) has been predicted for the ^{20}Ne isotope nuclei within the nuclear shell model, by using KUOSD interactions. With the harmonic oscillator potential (HO, b), the transition probabilities were calculated for $b < 0$ [13]. The comparison between experimental B (E2) and theoretical shows an advantage for KUOSD calculations for many levels. The primary polarization effect was included by selecting the effective charge for proton and neutrons $e_p = e_n = 0.350$ respectively. Comparison of calculated values with available experiment data [12] is shown in Table.2 ,New electric B(E2) transitions were expected in our results by using (KUOSD) interaction of the studied isotope. transition probabilities B(E2) gives agreeable agreement comparing with experimental data.

Table2.Comparison of the B(E2) results for ^{20}Ne nucleus with the experimental data[12]

$J_j^+ \rightarrow J_i^+$	B(E2) our.Results ($e^2 \text{ fm}^4$) $e_n=0.350, e_p=0.350$	B(E2) Exp.Results ($e^2 \text{ fm}^4$)
$2_1 \rightarrow 0_1$	59.04	65.46
$2_2 \rightarrow 0_1$	0.0098	2.35
$4_1 \rightarrow 2_1$	73.53	70.94
$2_2 \rightarrow 2_1$	2.612	5.48
$0_2 \rightarrow 2_1$	13.98	11.60
$3_1 \rightarrow 2_1$	0.0297	
$4_2 \rightarrow 2_1$	3.447	26.76
$1_1 \rightarrow 2_1$	0.3683	
$3_2 \rightarrow 2_1$	0.3011	
$1_2 \rightarrow 2_1$	0.2690	
$2_2 \rightarrow 4_1$	1.186	
$6_1 \rightarrow 4_1$	59.99	64.49
$3_1 \rightarrow 4_1$	0.5387	
$4_2 \rightarrow 4_1$	0.2108	
$3_2 \rightarrow 4_1$	0.1390	
$5_1 \rightarrow 4_1$	0.3331	
$5_2 \rightarrow 4_1$	0.8671	
$6_2 \rightarrow 4_1$	0.7837	
$2_2 \rightarrow 0_2$	18.76	
$4_2 \rightarrow 6_1$	10.49	
$8_1 \rightarrow 6_1$	38.35	29.02
$5_1 \rightarrow 6_1$	2.560	
$5_2 \rightarrow 6_1$	0.04356	
$6_2 \rightarrow 6_1$	0.9021	
$7_1 \rightarrow 6_1$	0.0006437	
$7_2 \rightarrow 6_1$	0.1091	
$8_2 \rightarrow 6_1$	1.464	
$3_1 \rightarrow 2_2$	10.31	
$4_2 \rightarrow 2_2$	0.1504	
$1_1 \rightarrow 2_2$	0.6821	
$3_2 \rightarrow 2_2$	0.8443	
$1_2 \rightarrow 2_2$	0.7005	
$1_1 \rightarrow 3_1$	20.57	
$3_2 \rightarrow 3_1$	8.108	
$1_2 \rightarrow 3_1$	0.04564	
$5_1 \rightarrow 3_1$	26.23	
$5_2 \rightarrow 3_1$	0.2754	
$1_1 \rightarrow 3_2$	28.82	
$1_2 \rightarrow 3_2$	0.007264	
$5_1 \rightarrow 3_2$	0.1324	
$5_2 \rightarrow 3_2$	11.96	
$3_1 \rightarrow 4_2$	3.177	
$3_2 \rightarrow 4_2$	5.354	
$5_1 \rightarrow 4_2$	7.665	

$5_2 \rightarrow 4_2$	13.90	
$6_2 \rightarrow 4_2$	0.07915	
$5_2 \rightarrow 5_1$	1.578	
$6_2 \rightarrow 5_1$	17.20	
$7_1 \rightarrow 5_1$	18.72	
$7_2 \rightarrow 5_1$	1.740	
$6_2 \rightarrow 8_1$	0.8867	
$7_1 \rightarrow 8_1$	4.244	
$7_2 \rightarrow 8_1$	0.4456	
$8_2 \rightarrow 8_1$	2.637	
$6_2 \rightarrow 5_2$	1.162	
$7_1 \rightarrow 5_2$	3.502	
$7_2 \rightarrow 5_2$	0.6213	
$7_1 \rightarrow 6_2$	3.252	
$7_2 \rightarrow 6_2$	0.02658	
$8_2 \rightarrow 6_2$	26.69	
$7_2 \rightarrow 7_1$	5.588	
$8_2 \rightarrow 7_1$	13.41	
$8_2 \rightarrow 7_2$	1.940	
$1_2 \rightarrow 1_1$	9.675	

IV. CONCLUSIONS

The calculations were performed using the OXBASH code. The study demonstrated that the interaction files are consistent with the experimental data available newly. Calculations of energy levels and reduced probabilities of the $B(E2)$ transitions for ^{20}Ne were carried out within the sd shell by using effective interaction KUOSD. We deduce the shell model configuration mixing are very successful in this model space. In our calculations, new energy levels have been predicted and many levels that are not experimentally confirmed have been confirmed.

REFERENCES

- [1] Mohammadi, S., S. Arbab, and E. Tavakoli. "Energy Levels Calculations of 22-23Na and 24-26Mg Isotopes Using Shell Model Code OXBASH." American Journal of Modern Physics. Special Issue: Many Particle Simulations 4.3-1 (2015): 27-31.
- [2] Povh, Bogdan, Klaus Rith, and Frank Zetsche. Particles and nuclei. Vol. 4. Berlin: Springer, 1995.
- [3] Brown BA, Arima A, McGrory JB (1977) E2 core-polarization charge for nuclei near 16O and 40Ca. Nuclear Physics 277.
- [4] Akiyama Y, Arima A, Sebe T (1969) The structure of the SD shell nuclei: (IV). 20Ne, 21Ne, 22Ne, 22Na and 24Mg. Nuclear Physics 138: 273-304
- [5] Radhi, A., G. N. Flaiyh, and E. M. Raheem. "Electric Quadrupole Transitions of Some Even-Even Neon Isotopes." Iraqi Journal of Science 56.2A (2015): 1047-1060.

- [6] B McGrory, a.B.H.W., Large-Scale Shell-Model Calculations. Annual Review of Nuclear and Particle Science, 1980. 30: p. 383.
- [7] M. Honma, , B. A. Brown, T. Mizusaki and T. Otsuka, Nucl. Phys. A704 (2002), 134c .
- [8] H. Eng " introduction to nuclear physics " Addison Wesley, 1966.
- [9] R. D. Lawson. Theory of the nuclear shell model. Oxford: Clarendon Press, 1980.
- [10] OXBASH for Windows, B. A. Brown, A. Etchegoyen, N. S.Godwin, W. D.M. Rae, W.A. Richter, W.E. Ormand, E.K.Warburton, J.S. Winfield, L. Zhao and C.H. Zimmerman,MSU_NSCL report number 1289.
- [11] Al-Sammarraie, Anwer A., Hasan Abu Kassim, and F. I. Sharrad. "Nuclear Structure for 24Mg within sd-Shell Model Space Hamiltonians." Armenian J. Phys. 8 (2015): 170-179.
- [12] <https://www.nndc.bnl.gov/nudat2/chartNuc.jsp>.
- [13] Hasan, A. K. "Shell Model Calculations for 18, 19, 200 Isotopes by Using USDA and USDB Interactions." Ukrainian Journal of Physics 63.3 (2018): 189-189.

AUTHORS

First Author – Ali k.Hasan, Department of Physics, College of Education for Girls, University of Kufa, Iraq, alikh.alsinayyid@uokufa.edu.iq
Second Author – Hadeel H.Abed, Department of Physics, College of Education for Girls, University of Kufa, Iraq, hadeelh.shamamrrah@student.uokufa.edu.iq